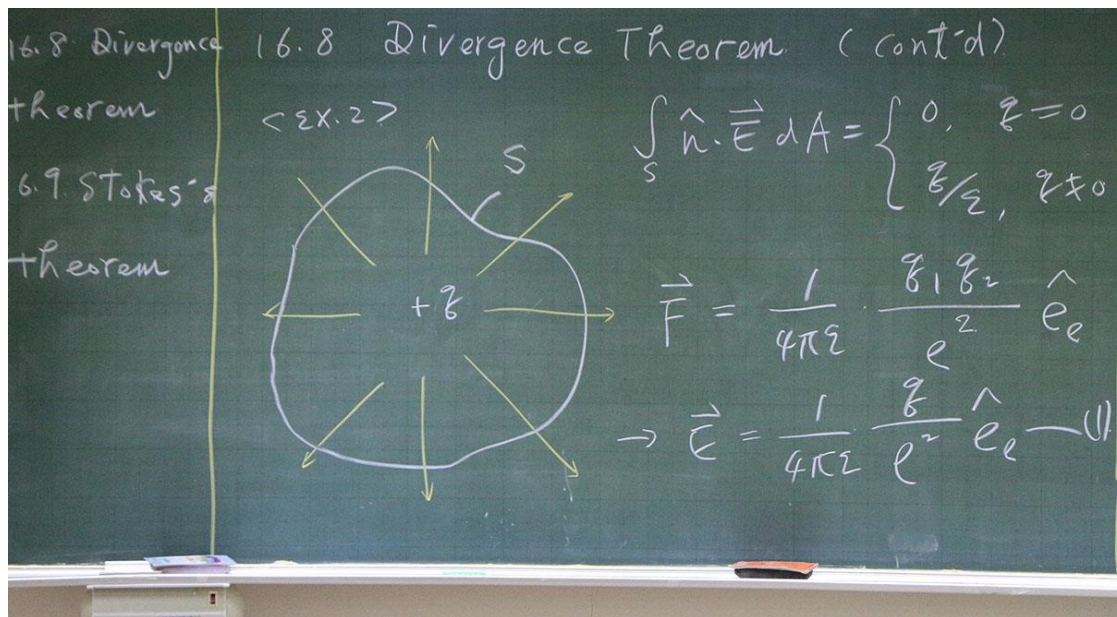
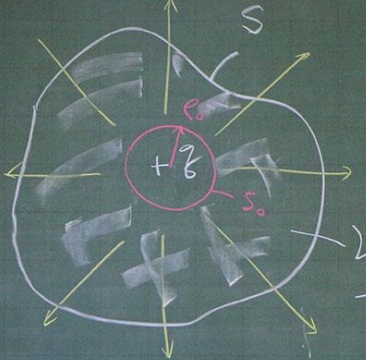




16.8 Divergence theorem
16.9 Stokes's theorem

16.8 Divergence Theorem (cont'd)

<ex. 2>



$\int_S \hat{n} \cdot \vec{E} dA = \begin{cases} 0, & q=0 \\ q/\epsilon, & q \neq 0 \end{cases}$

$\vec{F} = \frac{1}{4\pi\epsilon} \frac{q_1 q_2}{r^2} \hat{e}_r$

$\rightarrow \vec{E} = \frac{1}{4\pi\epsilon} \frac{q}{r^2} \hat{e}_r \quad (1)$

1. S does NOT enclose q (i.e., $q=0$)

$\int_S \hat{n} \cdot \vec{E} dA = \int_V \nabla \cdot \vec{E} dV = 0 \quad (1)$

2. S does enclose q ! ^{spherical coordinates.}

$\int_{V'} \nabla \cdot \vec{E} dV = 0 = \int_S \hat{n} \cdot \vec{E} dA + \int_{S_0} \hat{n} \cdot \vec{E} dA$

$$\int_S \hat{n} \cdot \vec{E} dA = - \int_{S_0} \hat{n} \cdot \vec{E} dA = - \int_{S_0} (-\hat{e}_e)$$

$$\left(\frac{1}{4\pi\epsilon} \frac{q}{\ell_0^2} \right) \hat{e}_e dA = \int_{S_0} \left(\frac{1}{4\pi\epsilon} \frac{q}{\ell_0^2} \right) dA$$

$$= \frac{q}{4\pi\epsilon \ell_0^2} \times 4\pi\ell_0^2 = \frac{q}{\epsilon} \quad \#$$

PS

$$\int_S \hat{n} \cdot \vec{E} dA = \int_V \nabla \cdot \vec{E} dV = \frac{q}{\epsilon}$$

$$= \int_V \frac{q_V}{\epsilon} dV \quad \therefore \text{integrand should}$$

$$\text{be the same} \rightarrow \nabla \cdot \vec{E} = \frac{q_V}{\epsilon} \quad \#$$

8. Divergence theorem
 9. Stokes' theorem

EX.3. Continuity eqn. of fluid mechanics

$M = \int_V \rho \, dV$

$$1. \frac{dM}{dt} = \frac{d}{dt} \int_V \rho(x, y, z, t) \, dV$$

$$= \int_V \frac{\partial \rho}{\partial t} \, dV \quad \text{--- (1)}$$

2. flow into V thru its surface

$$- \int \rho \hat{n} \cdot \vec{v} \, dA \quad \text{--- (2)}$$

$\frac{\text{kg}}{\text{m}^3} \cdot \frac{\text{m}}{\text{sec}} \cdot \text{m}^2 = \frac{\text{kg}}{\text{sec}}$

3. mass generated by chemical reactions

$$\int_V f(x, y, z, t) \, dV \quad \text{--- (3)}$$

(1) = (2) + (3), we obtain,

$$\int_V \frac{\partial \sigma}{\partial t} dV = - \int_S \sigma \hat{n} \cdot \vec{v} dA + \int_V f dV \quad (4)$$

$$+ \int_S \sigma \hat{n} \cdot \vec{v} dA = \int_S \hat{n} \cdot (\sigma \vec{v}) dA$$

$$= \int_V \nabla \cdot (\sigma \vec{v}) dV \quad (5)$$

Therefore, (4) becomes,

$$\int_V \left[\frac{\partial \sigma}{\partial t} + \nabla \cdot (\sigma \vec{v}) - f \right] dV = 0$$

$\therefore$ V is arbitrary $\therefore [] = 0$

$$\Rightarrow \frac{\partial \sigma}{\partial t} + \nabla \cdot (\sigma \vec{v}) - f = 0 \quad (6)$$

8. Divergence

theorem

9. Stokes's

theorem

PS

1. There is no chemical reaction,

$$f=0 \rightarrow \frac{\partial \sigma}{\partial t} + \nabla \cdot (\sigma \vec{v}) = 0$$

2. $f=0$, incompressible field

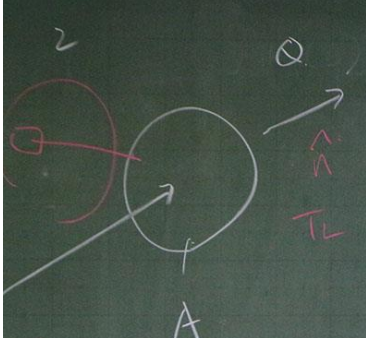
$$\sigma = \text{const.} \therefore \frac{\partial \sigma}{\partial t} = 0, \quad \nabla \cdot (\sigma \vec{v})$$

$$= \cancel{\nabla \sigma} \cdot \vec{v} + \sigma \nabla \cdot \vec{v} \rightarrow \nabla \cdot \vec{v} = 0$$

< ex. 4. Unsteady-state heat conduction >

$$\begin{aligned} 1. \quad H &= m c T \rightarrow \frac{dH}{dt} = \frac{d}{dt} \int_V c T \sigma dV \\ &= \int_V c \sigma \frac{\partial T}{\partial t} dV \quad \text{--- (1)} \end{aligned}$$

2



Calories/sec $\rightarrow -k \cdot A \frac{\partial T}{\partial n}$ Fourier conduction (2)

$\rightarrow + \int_S (k \frac{\partial T}{\partial n}) dA$

3. chemical reaction, mechanical hysteresis, f

$\int_V f(x, y, z, t) dV$

$\frac{\partial T}{\partial n} = \frac{(T_L - T_H)}{\Delta L}$

(1) = (2) + (3) $\rightarrow$ conservation of Energy !!

$\int_V c \rho \frac{\partial T}{\partial t} dV = \int_S \hat{n} \cdot (k \nabla T) dA$

$k \frac{\partial T}{\partial n} = k \cdot (\nabla T \cdot \hat{n})$

$+ \int_V f dV$ (4)

where, $\int_S \hat{n} \cdot (k \nabla T) dA = \int_V \nabla \cdot (k \nabla T) dV$

$$\int_V \left[c \sigma \frac{\partial T}{\partial t} - \nabla \cdot (k \nabla T) - f \right] dV = 0 \quad (5)$$

$$PS: \frac{\partial T}{\partial n} = \nabla T \cdot \hat{n}$$

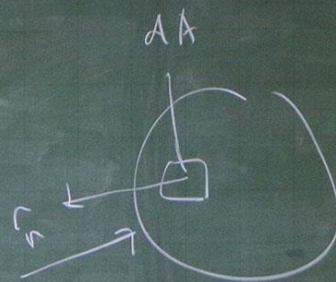
directional
derivative

$$c \sigma \frac{\partial T}{\partial t} - k \nabla^2 T - f = 0$$

$$1. f = 0$$

$$c \sigma \frac{\partial T}{\partial t} - k \nabla^2 T = 0$$

$$\left(\frac{k}{c \sigma} \right) \nabla^2 T = 1 \cdot \frac{\partial T}{\partial t} \Rightarrow \alpha^2 \nabla^2 T = \frac{\partial T}{\partial t} \quad \#$$



<EX. 5>

$$J = -D \frac{\partial C}{\partial n} \Rightarrow \beta^2 \nabla^2 C = \frac{\partial C}{\partial t} \quad \#$$

Fick's 2nd law

Divergence
rem

$$\alpha^2 \nabla^2 T = \frac{\partial T}{\partial t} \rightarrow 0 \rightarrow \text{steady state}$$

Takes a

$\neq 0 \rightarrow \text{unsteady state}$

rem

$$\rightarrow \nabla^2 T = 0 \quad \text{"Laplace Eqn."}$$

ch. 18: diffusion eqn

ch. 19: wave eqn

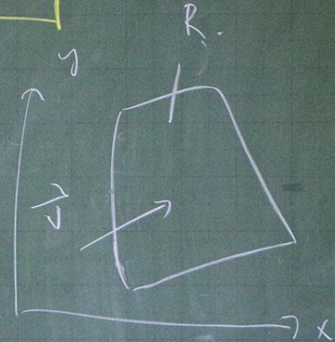
ch. 20: Laplace

$$c^2 \nabla^2 E = \frac{\partial^2 T}{\partial t^2}$$

< 2D case of Divergence theorem >

$$\int_V \nabla \cdot \vec{V} dV = \int_S \hat{n} \cdot \vec{V} dA$$

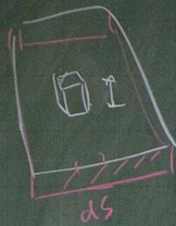
$$\vec{V} = V_x(x, y) \hat{i} + V_y(x, y) \hat{j}$$



sol.

$$dV = dA(1) = dA$$

$$dA = dS(1) = dS$$



LHS of (*)

$$\int_V \nabla \cdot \vec{V} dV = \int_V \left(\frac{\partial V_x}{\partial x} + \frac{\partial V_y}{\partial y} + 0 \right) dA$$

$$= \int_S \left(\frac{\partial V_x}{\partial x} + \frac{\partial V_y}{\partial y} \right) dA = \int_S \nabla \cdot \vec{V} dA \quad (1)$$

RHS of (*)

1. Top / bottom faces: $\hat{n} \cdot \vec{v} = 0$ ($\because \hat{n} = \pm \hat{k}$)

$$\rightarrow \hat{k} \cdot (V_x \hat{i} + V_y \hat{j}) = 0$$

2. for the other 4 faces,

$$\int_S \hat{n} \cdot \vec{v} dA = \oint_C \hat{n} \cdot \vec{v} dS$$

$$\int_S \nabla \cdot \vec{v} dA = \oint_C \hat{n} \cdot \vec{v} dS \quad \#$$

Divergence

theorem

Stokes's

theorem

16.9. Stokes's theorem.

I. Background

$$\left\{ \begin{array}{l} \int f dl = \int \vec{F} \cdot d\vec{R} \\ \int \vec{H} \cdot d\vec{R} = I \\ \int \vec{E} \cdot d\vec{R} \end{array} \right.$$

